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### 1.4. Stability

## A. Gain Objective

The principal aim of a feedback loop is to set an so that is a reverse $\beta_{\mathrm{FB}}$ translation of $\mathrm{s}_{\mathrm{I}}$ 's mirrored reflection. For $\mathrm{A}_{\mathrm{CL}}$ to follow this translation, $1 / \beta_{\mathrm{FB}}$ should be lower than $\mathrm{A}_{\mathrm{FW}}$. But since gain is another goal, $1 / \beta_{\mathrm{FB}}$ should be one or greater. So in practice, $A_{F W}$ is usually higher than $1 / \beta_{F B}$ across frequencies of interest and $\beta_{\mathrm{FB}}$ is lower than or equal to 1 or 0 dB .

## B. Stability Criterion

High $A_{L G}$ is desirable in feedback systems because amplifying $S_{E}$ reduces the mismatch between $\mathrm{S}_{\mathrm{I}}$ and $\mathrm{S}_{\mathrm{FB}}$. Translating $\mathrm{so}_{\mathrm{O}}$ to $\mathrm{S}_{\mathrm{FB}}$, comparing $\mathrm{S}_{\mathrm{FB}}$ to $\mathrm{S}_{\mathrm{I}}$, and amplifying the resulting $\mathrm{s}_{\mathrm{E}}$ so this $\mathrm{A}_{\mathrm{LG}}$ is high and $\mathrm{s}_{\mathrm{O}}$ is accurate usually requires two or more stages. Since each stage incorporates one or more poles, finding two or more poles in $\mathrm{A}_{\mathrm{LG}}$ is not uncommon.

In Fig. 5, to cite an example, $\mathrm{A}_{\mathrm{LG}}$ 's zero- or low-frequency gain $\mathrm{A}_{\mathrm{LG} 0}$ is well above 0 dB . $\mathrm{A}_{\mathrm{LG}}$ falls 20 dB per decade after $\mathrm{p}_{1}$ and another 20 dB per decade after $\mathrm{p}_{2}$. $\mathrm{A}_{\mathrm{LG}}$ crosses 0 dB at a unity-gain frequency $\mathrm{f}_{0 \mathrm{~dB}}$ that is higher than $\mathrm{p}_{1}$ and $\mathrm{p}_{2}$. Since each pole reduces phase shift up to $90^{\circ}, \angle \mathrm{A}_{\mathrm{LG}}$ reaches $-180^{\circ}$ (at the inversion frequency $\mathrm{f}_{180^{\circ}}$ ) before $\mathrm{A}_{\mathrm{LG}}$ crosses 0 dB .


Fig. 5. Unstable loop-gain response.
Since $A_{\text {LG }}$ inverts with $-180^{\circ}$ past $f_{180^{\circ}}, A_{L G}$ is -1 at $f_{0 d B}$. With this much phase shift, positive feedback peaks $A_{C L}$ at $f_{0 d B}$ towards infinity:

$$
\begin{equation*}
\mathrm{A}_{\mathrm{CL}}=\mathrm{A}_{\mathrm{FW}} \| \frac{1}{\beta_{\mathrm{FB}}}=\left.\frac{\mathrm{A}_{\mathrm{FW}}}{1+\mathrm{A}_{\mathrm{LG}}}\right|_{\mathrm{A}_{\mathrm{LG}}=1 \angle 180^{\circ}}=\frac{\mathrm{A}_{\mathrm{FW}}}{1-1} \rightarrow \infty \tag{8}
\end{equation*}
$$

Since $A_{V}$ 's $R_{I N}$ is very high, $\beta_{F B}$ is the $v_{0}$ fraction that $R_{2}$ sets across $R_{1}$ :

$$
\begin{equation*}
\beta_{\mathrm{FB}} \equiv \frac{\mathrm{~V}_{\mathrm{FB}}}{\mathrm{v}_{\mathrm{O}}} \approx \frac{\mathrm{R}_{1}}{\mathrm{R}_{1}+\mathrm{R}_{2}} . \tag{20}
\end{equation*}
$$

So $A_{L G}$ is $A_{F W} \beta_{F B}$ and $A_{L G}$ reaches 0 dB at $\mathrm{A}_{\mathrm{LG} 0} \mathrm{p}_{\mathrm{A}}$ or $\mathrm{A}_{\mathrm{FW} 0} \beta_{\mathrm{FB}} \mathrm{P}_{\mathrm{A}}$ :

$$
\begin{gather*}
\mathrm{A}_{\mathrm{LG}}=\mathrm{A}_{\mathrm{FW}} \beta_{\mathrm{FB}} \approx\left(\frac{\mathrm{~A}_{\mathrm{V} 0}}{1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{A}}}\right)\left(\frac{\mathrm{R}_{1}}{\mathrm{R}_{1}+\mathrm{R}_{2}}\right)  \tag{21}\\
\mathrm{f}_{\mathrm{odB}} \approx \mathrm{~A}_{\mathrm{LG} 0} \mathrm{p}_{\mathrm{A}}=\mathrm{A}_{\mathrm{FW} 0} \beta_{\mathrm{FB}} \mathrm{p}_{\mathrm{A}} \approx \mathrm{~A}_{\mathrm{V} 0}\left(\frac{\mathrm{R}_{1}}{\mathrm{R}_{1}+\mathrm{R}_{2}}\right) \mathrm{p}_{\mathrm{A}} . \tag{22}
\end{gather*}
$$

And the voltage gain $\mathrm{A}_{\mathrm{vo}}$ to vo is $\mathrm{A}_{\mathrm{CL}}$ 's $\mathrm{A}_{\mathrm{vo}}\| \|^{1 / \beta_{\mathrm{FB}}}$ up to $\mathrm{f}_{\mathrm{odB}}$ :

$$
\begin{equation*}
\mathrm{A}_{\mathrm{Vo}} \equiv \frac{\mathrm{v}_{\mathrm{O}}}{\mathrm{v}_{\mathrm{IN}}}=\mathrm{A}_{\mathrm{FW}} \| \frac{1}{\beta_{\mathrm{FB}}} \approx\left(\mathrm{~A}_{\mathrm{V} 0} \| \frac{\mathrm{R}_{1}+\mathrm{R}_{2}}{\mathrm{R}_{1}}\right)\left(\frac{1}{1+\mathrm{s} / 2 \pi \mathrm{f}_{\mathrm{odB}}}\right), \tag{23}
\end{equation*}
$$

which reduces to $1 / \beta_{\mathrm{FB}}{ }^{\prime} \mathrm{s}\left(\mathrm{R}_{1}+\mathrm{R}_{2}\right) / \mathrm{R}_{1}$ up to $\mathrm{f}_{\text {odB }}$ when $\mathrm{A}_{\mathrm{FW}}$ 's $\mathrm{A}_{\mathrm{Vo}}$ is much greater than this $1 / \beta_{\mathrm{FB}}$.


Fig. 15. Non-inverting (voltage-mixed) op amp.

Example 1: Determine $\mathrm{A}_{\mathrm{FW} 0}, \beta_{\mathrm{FB}}, \mathrm{A}_{\mathrm{LG} 0}, \mathrm{f}_{\mathrm{odB}}, \mathrm{A}_{\mathrm{voo}}$, and $\mathrm{f}_{\mathrm{CL}(\mathrm{BW})}$ when $\mathrm{A}_{\mathrm{vo}}$ is $100 \mathrm{~V} / \mathrm{V}, \mathrm{p}_{\mathrm{A}}$ is $10 \mathrm{kHz}, \mathrm{R}_{1}$ is $10 \mathrm{k} \Omega$, and $\mathrm{R}_{2}$ is $90 \mathrm{k} \Omega$.

## Solution:

$$
\begin{aligned}
& \mathrm{A}_{\mathrm{FW} 0} \approx \mathrm{~A}_{\mathrm{V} 0}=100 \mathrm{~V} / \mathrm{V} \\
& \beta_{\mathrm{FB}} \approx \frac{\mathrm{R}_{1}}{\mathrm{R}_{1}+\mathrm{R}_{2}}=\frac{10 \mathrm{k}}{10 \mathrm{k}+90 \mathrm{k}}=100 \mathrm{mV} / \mathrm{V} \\
& \mathrm{~A}_{\mathrm{LG} 0}=\mathrm{A}_{\mathrm{FW} 0} \beta_{\mathrm{FB}} \approx(100)(100 \mathrm{~m})=10 \mathrm{~V} / \mathrm{V} \\
& \mathrm{f}_{0 \mathrm{~dB}} \approx \mathrm{~A}_{\mathrm{LG} 0} \mathrm{p}_{\mathrm{A}} \approx(10)(10 \mathrm{k})=100 \mathrm{kHz}
\end{aligned}
$$

phase from shifting $180^{\circ}$. This way, $\mathrm{A}_{\mathrm{LG}}$ follows $\mathrm{A}_{\mathrm{S}}$ up to $\mathrm{p}_{1}$ and continues to fall after $z_{S 1}$ and $z_{S 2}$ in $A_{S}$ counter the effects of $p_{1}$ and $p_{2}$ in $A_{L G}$.

Parasitic poles in $\mathrm{A}_{\mathrm{S}}$ eventually limit $\mathrm{A}_{\mathrm{S}}$ 's bandwidth. So after $\mathrm{Z}_{\mathrm{S} 1}$ and $\mathrm{z}_{\mathrm{S} 2}, \mathrm{~A}_{\mathrm{S}}$ flattens with $\mathrm{p}_{\mathrm{S} 2}$ and falls with $\mathrm{p}_{\mathrm{S} 3}$. Although $\mathrm{p}_{\mathrm{S} 2}$ and $\mathrm{p}_{\mathrm{S} 3}$ are not always apart, only one of these poles can be close to $f_{0 \mathrm{~dB}}$ for stability.

### 3.2. Amplifier Translations

An op amp can add $\mathrm{p}_{\mathrm{s} 1}$. This op amp, however, cannot be any op amp. This is because the low-frequency gain $\mathrm{A}_{\mathrm{S} 0}$ and $\mathrm{p}_{\mathrm{S} 1}$ that $\mathrm{A}_{\mathrm{V} 0}$ and $\mathrm{p}_{\mathrm{A}}$ set should establish an $f_{0 \mathrm{~dB}}$ that keeps the feedback system stable:

$$
\begin{equation*}
\mathrm{A}_{\mathrm{S}} \approx \frac{\mathrm{~A}_{\mathrm{v} 0}}{1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{A}}} \tag{36}
\end{equation*}
$$

The OTAs in Fig. 22 can also add $\mathrm{p}_{\mathrm{S} 1} . \mathrm{A}_{\mathrm{S} 0}$ is the gain that $\mathrm{A}_{\mathrm{G}}$ sets across $R_{F}$. In the first implementation, $A_{S}$ falls past $p_{F}$ when $C_{F}$ shunts $R_{F}$ :

$$
\begin{equation*}
A_{S} \approx A_{G}\left(R_{F} \| \frac{1}{S_{F}}\right)=\frac{A_{G} R_{F}}{1+\mathrm{sR}_{F} C_{F}}=\frac{\mathrm{A}_{\mathrm{G}} \mathrm{R}_{\mathrm{F}}}{1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{F}}} \tag{37}
\end{equation*}
$$



Fig. 22. Dominant-pole and pole-zero OTAs.
Current-limiting $\mathrm{C}_{\mathrm{F}}$ with $\mathrm{R}_{\mathrm{C}}$ adds $\mathrm{Z}_{\mathrm{S} 1}$. With $\mathrm{R}_{\mathrm{C}}, \mathrm{A}_{\mathrm{S}}$ falls past $\mathrm{p}_{\mathrm{C}}$ when $C_{F}$ shunts $R_{C}$ and $R_{F}$ before parasitic capacitance $C_{X}$ at vo shunts $R_{F}$. $p_{C}$ eventually fades past $z_{C X}$ when $C_{F}$ shorts with respect to $\mathrm{R}_{\mathrm{C}}$. Once shorted, As flattens to $A_{G}\left(R_{F} \| R_{C}\right)$ and later falls past po when $C_{X}$ shunts $R_{F} \| R_{C}$ :

$$
\begin{aligned}
A_{S} & =A_{G}\left[R_{F}\left\|\left(Z_{F}+R_{C}\right)\right\| Z_{X}\right] \\
& =\frac{A_{G} R_{F}\left(1+\mathrm{sC}_{F} R_{C}\right)}{\mathrm{s}^{2} R_{C} C_{F} R_{F} C_{X}+\mathrm{s}\left[\left(\mathrm{R}_{F}+R_{C}\right) C_{F}+R_{F} C_{X}\right]+1}
\end{aligned}
$$



Fig. 31. Dominant-pole inverting mixed translation.
$\mathrm{A}_{\mathrm{S}}$ follows $\mathrm{A}_{\mathrm{F}}$ 's $\mathrm{A}_{\mathrm{V} 0}$ until $\mathrm{A}_{\beta}$ drops below $\mathrm{A}_{\mathrm{V} 0}$ at $\mathrm{p}_{\mathrm{X} 1}$. With two poles in $A_{F}$ and one in $A_{\beta}, A_{F}$ falls faster than $A_{\beta}$. As a result, $A_{S}$ falls with $A_{\beta}$ past $\mathrm{p}_{\mathrm{X} 1}$ until $\mathrm{A}_{\mathrm{F}}$ falls below $\mathrm{A}_{\beta}$ at $\mathrm{p}_{\mathrm{X} 2}$. This way, $\mathrm{A}_{\mathrm{S} 0}$ is $-\mathrm{A}_{\mathrm{V} 0}, \mathrm{p}_{\mathrm{S} 1}$ is $\mathrm{p}_{\mathrm{X} 1}$, and $p_{S 2}$ is $p_{\mathrm{X} 2}$, but only when $\mathrm{A}_{\beta}$ 's projection to $\mathrm{p}_{\mathrm{X} 1}$ precedes $\mathrm{p}_{\mathrm{A}}$ and $\mathrm{p}_{\mathrm{F}}$ and $\mathrm{A}_{\mathrm{F}}$ 's projection to $\mathrm{p}_{\mathrm{X} 2}$ exceeds $\mathrm{p}_{\mathrm{X} 1}$ :

$$
\begin{gather*}
\left.\left|\mathrm{A}_{\beta}\right| \approx \frac{1}{\mathrm{SR}_{\mathrm{F}} \mathrm{C}_{\mathrm{F}}}\right|_{\mathrm{f}_{\mathrm{o}} \geq \frac{1}{2 \pi \mathrm{~A}_{\mathrm{V} 0} \mathrm{R}_{\mathrm{F}} \mathrm{C}_{\mathrm{F}}}=\frac{\mathrm{p}_{\mathrm{F}}}{\mathrm{~A}_{\mathrm{V} 0}} \approx \mathrm{p}_{\mathrm{X} 1}} \leq\left|\mathrm{A}_{\mathrm{F}}\right|_{\mathrm{f}_{0}<\mathrm{p}_{\mathrm{A}}} \approx \mathrm{~A}_{\mathrm{V} 0},  \tag{65}\\
\left.\left|\mathrm{~A}_{\mathrm{F}}\right|_{\mathrm{f}_{\mathrm{O}}>\mathrm{p}_{\mathrm{A}}, \mathrm{p}_{\mathrm{F}}} \approx \frac{\mathrm{~A}_{\mathrm{V} 0} \mathrm{p}_{\mathrm{A}} \mathrm{p}_{\mathrm{F}}}{\mathrm{f}_{\mathrm{O}}^{2}}\right|_{\mathrm{f}_{\mathrm{o}} \geq \mathrm{A}_{\mathrm{V} 0} \mathrm{p}_{\mathrm{A}} \approx \mathrm{p}_{\mathrm{X} 2}} \leq\left|\mathrm{A}_{\beta}\right|=\frac{\mathrm{p}_{\mathrm{F}}}{\mathrm{f}_{\mathrm{O}}}, \tag{66}
\end{gather*}
$$

$$
\mathrm{A}_{\mathrm{S}}=\mathrm{A}_{\mathrm{F}} \| \mathrm{A}_{\beta} \approx \frac{-\mathrm{A}_{\mathrm{v} 0}}{\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{x} 1}\right)\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{x} 2}\right)}
$$



Fig. 32. Pole-zero inverting mixed translation.
Current-limiting $C_{F}$ with $R_{C}$ in Fig. 32 reverses $C_{F}$ 's pole in $A_{F}$ and $A_{\beta}$. So $A_{F}$ starts with $-A_{V 0}, A_{F}$ falls past $A_{V}$ 's $p_{A}$ and $p_{C}$ when $C_{F}$ shunts $R_{C}$ and $\mathrm{R}_{\mathrm{F}}$, and $\mathrm{z}_{\mathrm{CX}}$ reverses $\mathrm{p}_{\mathrm{C}}$ when $\mathrm{R}_{\mathrm{C}}$ current-limits $\mathrm{C}_{\mathrm{F}}$. $\mathrm{A}_{\beta}$ falls as $\mathrm{C}_{\mathrm{F}}$ shorts and flattens to $-R_{C} / R_{F}$ past $z_{C X}$ when $C_{F}$ shorts with respect to $R_{C}$ :

$$
\begin{equation*}
A_{F} \approx\left(\frac{R_{C}+Z_{C}}{R_{F}+R_{C}+Z_{C}}\right)\left(-A_{V}\right)=\frac{-A_{V 0}\left(1+\mathrm{s} / 2 \pi \mathrm{z}_{\mathrm{CX}}\right)}{\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{A}}\right)\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{C}}\right)} \tag{68}
\end{equation*}
$$

energize and drain $\mathrm{L}_{\mathrm{X}}$. The duty-cycled inductance $\mathrm{L}_{\mathrm{DO}}$ is a $\mathrm{d}_{\mathrm{DO}}$ translation of $L_{X}$ with an $R_{L} / D_{D O}$ that is usually negligible in light of $R_{L D}$. So the static components of $d_{D O}, v_{E}$, and $v_{D}$ set $L_{D O}$ in Fig. 36 to $L_{X} / D_{D O}{ }^{2}$ and $A_{S L}$ to

$$
\begin{equation*}
\mathrm{A}_{\mathrm{SL}(C C M)} \equiv \frac{\mathrm{v}_{\mathrm{o}}}{\mathrm{~d}_{\mathrm{e}}{ }^{\prime}} \approx \frac{\left(\mathrm{V}_{\mathrm{E}}+\mathrm{V}_{\mathrm{D}}\right)\left(1+\mathrm{s} / 2 \pi \mathrm{z}_{\mathrm{C}}\right)\left(1-\mathrm{s} / 2 \pi \mathrm{z}_{\mathrm{DO}}\right)}{\mathrm{D}_{\mathrm{DO}}\left[\left(\mathrm{~s} / 2 \pi \mathrm{p}_{\mathrm{LC}}\right)^{2}+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{LC}} \mathrm{Q}_{\mathrm{LC}}+1\right]\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{SW}}\right)} . \tag{81}
\end{equation*}
$$



Fig. 36. Small-signal model of the switched inductor in CCM.
This gain drops as $L_{x}$ opens with frequency because $L_{X}$ feeds vo less current. A AL also falls as $\mathrm{C}_{\mathrm{O}}$ shorts and steers current away from $\mathrm{v}_{\mathrm{O}}$. The resulting inductor and capacitor poles $\mathrm{p}_{\mathrm{L}}$ and $\mathrm{p}_{\mathrm{C}}$ appear together as a double pole $\mathrm{p}_{\mathrm{LC}}$ at the transitional LC frequency $\mathrm{f}_{\mathrm{LC}}$ when $\mathrm{L}_{\mathrm{DO}}$ 's impedance $\mathrm{sL}_{\mathrm{DO}}$ overcomes $\mathrm{Co}^{\prime}$ 's $1 / \mathrm{sC} \mathrm{C}_{\mathrm{o}} . \mathrm{p}_{\mathrm{C}}$ eventually fades past $\mathrm{z}_{\mathrm{C}}$ when the capacitor resistance $\mathrm{R}_{\mathrm{C}}$ current-limits $\mathrm{Co}_{\mathrm{o}}$.

Duty-cycled outputs connect $\mathrm{L} x_{x}$ to vo only when draining $\mathrm{Lx}_{\mathrm{x}}$. So when the switching frequency $\mathrm{f}_{\mathrm{Sw}}$ is constant, extending $\mathrm{t}_{\mathrm{E}}$ shortens $\mathrm{Lx}^{\prime} \mathrm{s}$ drain time $t_{\mathrm{D}}$. Reducing drain current this way produces an inverting (out-ofphase) zero when the loss outpaces the gain. This duty-cycled zero $\mathrm{Z}_{\mathrm{DO}}$ normally appears above $\mathrm{p}_{\mathrm{LC}}$, but not by far. When present, $\mathrm{z}_{\mathrm{DO}}$ is usually below psw.
$\mathrm{p}_{\mathrm{LC}}$ is challenging because it shifts phase $180^{\circ}$ and peaks the gain. Since $\mathrm{L}_{\mathrm{DO}}$ 's and $\mathrm{Co}_{\mathrm{o}}$ 's impedances cancel at $\mathrm{f}_{\mathrm{LC}}$, inductor resistance $\mathrm{R}_{\mathrm{L}}$ and $\mathrm{R}_{\mathrm{C}}$ impose a series resistance $\mathrm{R}_{\mathrm{S}}$ that current-limits this peak. $\mathrm{R}_{\mathrm{LD}}$ dampens it below this level because $R_{L D}$ adds to the resistance that limits the $L C$ current. But since $R_{L}$ and $R_{C}$ are usually low and $R_{L D}$ is variable,

### 4.4. Current Mode

One way of eliminating $p_{L C}$ is by regulating $i_{\mathrm{L}}$. This way, the feedback translation that determines $i_{\mathrm{L}}$ is largely independent of $\mathrm{sL}_{\mathrm{x}}$. Removing this dependence to $\mathrm{sL}_{\mathrm{X}}$ eliminates the LC interaction that produces $\mathrm{p}_{\mathrm{LC}}$.

## A. Current Loop

$\mathrm{A}_{\text {IE }}$, the PWM, the switched inductor, and $\beta_{\text {IFB }}$ in Fig. 39 close an inverting feedback loop that sets $i_{\mathrm{L}}$. $\mathrm{A}_{\text {IE }}$ senses and amplifies the error that adjusts $d_{E}$ and $i_{L}$ so $v_{\text {IFb }}$ nears $v_{\text {Eo }}$. This way, $i_{L}$ is a reverse $\beta_{\text {Ifs }}$ translation of $v_{\text {Eo' }} s$ mirrored reflection, which is independent of $L_{x}$ 's impedance $\mathrm{sL}_{\mathrm{X}}$ :

$$
\begin{equation*}
\mathrm{i}_{\mathrm{L}}=\frac{\mathrm{v}_{\mathrm{IFB}}}{\beta_{\mathrm{IFB}}} \approx \frac{\mathrm{v}_{\mathrm{EO}}}{\beta_{\mathrm{IFB}}} . \tag{85}
\end{equation*}
$$

This is like removing $\mathrm{L}_{\mathrm{X}}$ from the circuit.


Fig. 39. Current-mode voltage controller.

## B. Loop Gain

When the forward gain $\mathrm{A}_{\mathrm{IF}}$ surpasses the feedback translation $\mathrm{A}_{\mathrm{I} \beta}$, the gain $\mathrm{A}_{\mathrm{G}}$ to $\mathrm{i}_{\mathrm{L}}$ follows $\mathrm{A}_{\mathrm{I}}$ 's $1 / \beta_{\mathrm{IFB}}$ up to the $\mathrm{p}_{\mathrm{G}}$ that the loop's $\mathrm{f}_{\text {IodB }}$ sets:

$$
\begin{equation*}
A_{G} \equiv \frac{i_{L}}{v_{\mathrm{EO}}}=\mathrm{A}_{\mathrm{IF}} \| \mathrm{A}_{\mathrm{I}} \approx \frac{1 / \beta_{\mathrm{IFB}}}{\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{G}}\right)\left(1+\mathrm{s} / 2 \pi \mathrm{p}_{\mathrm{SW}}\right)} . \tag{8}
\end{equation*}
$$

$A_{G}$ drops faster past $p_{\text {Sw }}$ when $f_{0}$ surpasses $f_{\text {Sw }}$. This $\beta_{\text {IFB }}$ is usually constant. So the loop that sets $i_{\mathrm{L}}$ in Fig. 39 is basically a bandwidth-limited transconductor that $\mathrm{d}_{\mathrm{DO}}$ in Fig. 40 duty-cycles.
$A_{L G}$ is the gain across $\beta_{F B}, A_{E}, A_{G}$, and $d_{D O}$ into $C_{o}$ with $R_{C}$ and $R_{L D}$. $\mathrm{A}_{\mathrm{LG}}$ starts with $\mathrm{A}_{\mathrm{E} 0} \mathrm{~A}_{\mathrm{G} 0} \mathrm{D}_{\mathrm{DO}} \mathrm{R}_{\mathrm{LD}} \beta_{\mathrm{FB}}$. $\mathrm{A}_{\mathrm{LG}}$ falls past $\mathrm{p}_{\mathrm{G}}, \mathrm{p}_{\mathrm{CP}}$, and $\mathrm{p}_{\mathrm{SW}}$ when
$\mathrm{A}_{\text {IF }} \mathrm{A}_{\text {PWM }} \mathrm{A}_{\text {IL }}$ is the part of $\mathrm{A}_{\text {ILG }}$ that determines feedback accuracy. This is because $\mathrm{A}_{\mathrm{G}}$ follows $\mathrm{A}_{I \beta}$ to the extent $\mathrm{A}_{I F}$ 's $\mathrm{A}_{\mathrm{IE}} \mathrm{A}_{\text {PWM }} \mathrm{A}_{I L}$ exceeds $\mathrm{A}_{I \beta}$, which is to say, $\mathrm{A}_{\mathrm{G}}$ approaches $1 / \beta_{\text {IFB }}$ when $\mathrm{A}_{\text {IF }}$ increases. In other words, regulation accuracy scales with $\mathrm{A}_{\mathrm{IF}}$.

### 5.3. Inherent Stability

As a stabilizer, the aim of $\mathrm{A}_{\text {IE }}$ is to ensure $\mathrm{A}_{\text {ILG }}$ reaches $\mathrm{f}_{\text {IodB }}$ with less than $180^{\circ}$ of phase shift. But since $A_{I L} ' s z_{C P}$ already recovers $90^{\circ}$ of the $180^{\circ}$ that $\mathrm{p}_{\mathrm{LC}}$ loses, $\mathrm{A}_{\mathrm{IE}}$ 's role can be to increase gain, and that way, extend $\mathrm{f}_{\mathrm{IOdB}}$. But for $\mathrm{f}_{\text {IodB }}$ to add no more than one pole $\mathrm{p}_{\mathrm{G}}$, $\mathrm{f}_{\text {IodB }}$ should be a decade or more below $\mathrm{A}_{\mathrm{IE}}$ 's bandwidth $\mathrm{p}_{\mathrm{IE} 1}$ and $\mathrm{f}_{\mathrm{Sw}}$ :

$$
\begin{equation*}
\left.\left.\mathrm{A}_{\mathrm{ILG}}\right|_{\mathrm{f}_{\mathrm{O}}>\mathrm{p}_{\mathrm{LC}}} \approx \frac{\mathrm{~A}_{\mathrm{ILG} 0} \mathrm{p}_{\mathrm{LC}}{ }^{2}}{\mathrm{z}_{\mathrm{CP}} \mathrm{f}_{\mathrm{O}}}\right|_{\mathrm{f}_{\mathrm{O}}=\mathrm{A}_{\mathrm{ILG} 0}\left(\frac{p_{\mathrm{LC}}{ }^{2}}{z_{\mathrm{CP}}}\right) \approx \mathrm{f}_{\mathrm{IOdB}}=\mathrm{p}_{\mathrm{G}} \leq \frac{\mathrm{p}_{\mathrm{IEI}}}{10}, \frac{\mathrm{f}_{\mathrm{SW}}}{10}}=1 . \tag{96}
\end{equation*}
$$

Since $A_{\text {ILG }}$ rises and falls to $0 \mathrm{~dB}, \mathrm{~A}_{\text {ILG }}$ usually starts low, which means $\mathrm{A}_{\text {IFo }}$ is also low. So $\mathrm{A}_{\text {IF }}$ in Fig. 44 starts low, climbs past $\mathrm{z}_{\mathrm{CP}}$, falls past pLC, and falls faster past psw. Although not always, $\mathrm{A}_{\text {IFo }}$ 's $\mathrm{A}_{\text {IE0 }} \mathrm{A}_{\text {PWm }} \mathrm{A}_{\text {IL0 }}$ is often lower than $A_{I \beta}$ 's $1 / \beta_{\text {IFB. }}$. So $A_{G}$ often starts with $\mathrm{A}_{\mathrm{IF})}$.


Fig. 44. Inherent transconductance in CCM.
$\mathrm{A}_{\mathrm{G}}$ climbs with $\mathrm{A}_{\text {IF }}$ past $\mathrm{z}_{\mathrm{CP}}$ until $\mathrm{A}_{\mathrm{IF}}$ surpasses $\mathrm{A}_{\mathrm{I}}$. This $\mathrm{z}_{\mathrm{CP}}$ is usually low because $C_{0}$ is high and $R_{L D}$ is moderate. Since $A_{I F}$ is the part of $A_{I L G}$ that excludes $\beta_{\text {IFB }}, \mathrm{A}_{\mathrm{ILG} 0}$ is below 1 when $\mathrm{A}_{\text {IF } 0}$ surpasses $1 / \beta_{\mathrm{IFB}}$. So $\mathrm{A}_{\mathrm{IF}}$ crosses $A_{I \beta}$ at a $p_{X 1}$ that is $1 / \mathrm{A}_{\mathrm{ILG} 0}$ times greater than $\mathrm{Z}_{\mathrm{CP}}$ :

$$
\begin{equation*}
\left.\mathrm{A}_{\mathrm{IF}}\right|_{\mathrm{f}_{\mathrm{O}}<\mathrm{P}_{\mathrm{LC}}}=\left.\mathrm{A}_{\mathrm{IE} 0} \mathrm{~A}_{\mathrm{PWM0}} \mathrm{~A}_{\mathrm{IL} 0}\left(\frac{\mathrm{f}_{\mathrm{O}}}{\mathrm{z}_{\mathrm{CP}}}\right)\right|_{\mathrm{f}_{0} \geq \frac{z_{\mathrm{CP}}}{\mathrm{~A}_{\mathrm{HLG}}} \approx \mathrm{p}_{\mathrm{XI} 1}>\mathrm{ZCP}} \geq \mathrm{A}_{\mathrm{IP}} \approx \frac{1}{\beta_{\mathrm{IFB}}} . \tag{97}
\end{equation*}
$$

$$
\begin{aligned}
& \mathrm{f}_{\mathrm{I} 0 \mathrm{~dB}} \approx \mathrm{~A}_{\mathrm{ILG} 0}\left(\frac{\mathrm{p}_{\mathrm{CS}}}{\mathrm{z}_{\mathrm{CP}}}\right) \mathrm{p}_{\mathrm{IE} 1}=(100)\left(\frac{120}{64}\right) \mathrm{p}_{\mathrm{IE} 1} \equiv \frac{\mathrm{f}_{\mathrm{SW}}}{10}=100 \mathrm{kHz} \\
& \therefore \quad \mathrm{p}_{\mathrm{IE} 1}=530 \mathrm{~Hz} \quad \text { and } \quad \mathrm{p}_{\mathrm{IE} 2} \geq \mathrm{f}_{\mathrm{I} 0 \mathrm{~dB}}=100 \mathrm{kHz} \\
& \mathrm{~A}_{\mathrm{G} 0}=\left(\mathrm{A}_{\mathrm{IE} 0} \mathrm{~A}_{\mathrm{PWM} 0} \mathrm{~A}_{\mathrm{IL} 0}\right) \| \frac{1}{\beta_{\mathrm{IFB}}} \\
& \quad \approx[(190)(2)(140 \mathrm{~m})] \| 1=980 \mathrm{~mA} / \mathrm{V}
\end{aligned}
$$

## 6. Digital Control

Feedback controllers use the voltage or current they sense to generate a pulsing command. From this perspective, feedback controllers are analogdigital converters (ADC). Mostly analog controllers mix, amplify, and stabilize the feedback system in the analog domain and mostly digital controllers in the digital domain.

Conventional ADCs digitize the voltage or current that digital controllers sense. Clocked digital-signal processors (DSP) use this digital word to mix, amplify, stabilize, and drive the switched inductor. Like analog controllers, digital controllers set loop gains that reach 0 dB with less than $180^{\circ}$ of phase, if possible, at the highest manageable $\mathrm{f}_{0 \mathrm{~dB}}$.


Fig. 47. Digital voltage-mode voltage controller.

### 6.1. Voltage Controller

Voltage-mode voltage controllers translate $\mathrm{v}_{\mathrm{O}}$ in Fig. 47 to $\mathrm{v}_{\mathrm{FB}}$ with $\beta_{\mathrm{FB}}$ and $\mathrm{v}_{\mathrm{FB}}$ into an N -bit digital word $\mathrm{d}_{1-\mathrm{N}}$ with ADCs. DSPs mix and compare this word $d_{1-\mathrm{N}}$ with a reference word $\mathrm{d}_{\mathrm{R}}$ and use the difference to output the pulsing command $\mathrm{d}_{\mathrm{E}}$ that adjusts $\mathrm{i}_{\mathrm{L}}$. This way, DSPs sense and amplify

