

# A Self-Synchronized Maximum-Power-Point Inductively Coupled Wireless Battery Charger for Embedded Microsensors

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**Abstract**—Embedded microsensors are critical in today’s fast-growing internet of things (IoT), as they provide an interface between the physical and digital worlds. An inductively coupled power receiver can replenish the onboard batteries and extend the microsensor’s lifetime. Among all power receivers, the switched resonant half-bridge receiver outputs the highest power as it can operate beyond the circuit’s breakdown adjust power transfer for maximum power point (MPP). However, switch synchronization and practical power transfer control remain two challenges. This paper presents a self-synchronization technique and a maximum-power-point adjustment control scheme for the switched resonant half-bridge wireless charger. The charger senses the high oscillation voltage without limiting the breakdown. Moreover, it can adjust the energy transfer frequency to reach and stay at the MPP when coupling varies. A prototype wireless charger is fabricated with 180 nm CMOS technology. The prototype charger, operating at 110 kHz, receives power up to 38 mW with a pair of  $11.7 \times 3.5 \times 2.6 \text{ mm}^3$  coils. Measurements show the charger outputs up to 89% of the available power across 0.067%–7.9% coupling range. The output power (in percentage of available power) and coupling range are  $1.3\times$  and  $13\times$  higher than the comparable state of the arts.

**Index Terms**—Internet of things, embedded microsensors, inductively coupled, wireless charger, maximum output power.

## I. POWERING MICROSENSORS INDUCTIVELY

Embedded microsensors are critical in the biomedical field and IoT [1-3]. However, these embedded microsensors’ tiny onboard batteries often drain quickly. Harvesting ambient energy, such as light or motion, can help replenish the battery. However, such energy sources are not always available. Often, the only option left is to recharge the battery wirelessly using a pair of inductively coupled coils.

Fig. 1 illustrates a typical inductively powered microsystem. To transfer power, the transmitter coil  $L_T$  in Fig. 1 runs an AC current and generates a changing magnetic field in the nearby space. An adjacent receiver coil  $L_R$  captures the magnetic flux that  $L_T$  emits and couples an electromotive force (EMF) voltage  $v_E$ . From the coupled  $v_E$ , the wireless charger draws power to charge up the energy storage  $v_B$ , which supplies the microsystem’s components, such as sensors, amplifiers, DSPs, etc. Designs in [3-6] combine the charging and supply stages into a single rectifying-regulating stage and remove the

energy storage  $v_B$ . The single-stage rectifying-regulating stage is, in essence, a supply, as it provides the power that the load requires. However, most embedded microsensors are heavily duty-cycled to conserve energy [7]. Therefore, the microsensors’ peak to average power ratio is often high [7]. Eliminating the intermediate energy storage  $v_B$  means that the transmitter needs to instantaneously supply the peak power that the load requires, so the required peak transmitting power is high. However, the transmitting power for many biomedical implant applications is often very limited due to safety and health standards [8]. For such single-stage rectifying-regulating supply systems, the design goal is to maximize the power conversion efficiency (PCE). As the voltage and power are regulated at the output, maximizing the receiver’s PCE minimizes the loss and conserves the most energy.

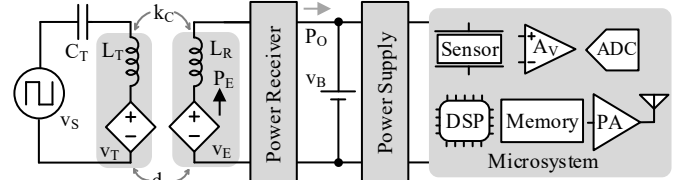


Fig. 1. Inductively powered microsystem.

The two-stage receiver and supply wireless power system in Fig. 1 works better for embedded microsensor applications. When the microsensor’s load idles, the inductively coupled power receiver charges the intermediate energy storage  $v_B$ . When the load power peaks, like a cushion, the energy stored in  $v_B$  helps supply the load, so the transmitter does not need to instantaneously supply the full peak power. As a result, the required peak transmitting power is lower, which is more friendly for health and safety. The wireless power receiver in Fig. 1 is, in essence, a charger, as it extracts power from  $v_E$  to fill up the energy storage  $v_B$ . The design goal of the wireless charger is to maximize the output power to charge up the battery as fast as possible. Since the coupling between the coils is often very low in embedded microsensor applications [9], the power receiver barely loads the transmitter. So the transmitter power remains about the same. Maximizing the charger’s output power equals maximizing the end-to-end efficiency.

The key to increasing the charger’s drawn power is to apply a high alternating voltage across the receiving coil [10]. The high coil voltage boosts the current from the coil so the receiver can draw more power from  $v_E$ . For MPP, the applied voltage is raised to the breakdown or the conduction loss limit.

The resonant bridge/half-bridge charger uses LC resonance to boost  $L_R$ ’s voltage [11-14]. However, since the

resonant capacitor and inductor are connected in parallel with the charger circuit. The maximum  $L_R$  voltage is limited by the charger circuit's breakdown. For standard deep-submicron CMOS circuits, this breakdown voltage is often low. The switched bridge [15-17] uses the rectified voltage to boost  $L_R$ 's voltage. Similarly, the  $L_R$ 's voltage is limited to the charger circuit's breakdown, as the rectified voltage connects with the circuit in parallel. The switched resonant half-bridge also uses LC resonance to boost  $L_R$ 's voltage. However, the resonant capacitor  $C_R$  in Fig. 2, connected in series with the charger circuit, bears most of the  $L_R$ 's voltage, so the breakdown limit is high. Therefore, the switched resonant half-bridge outputs the highest power without a buffer stage [10].

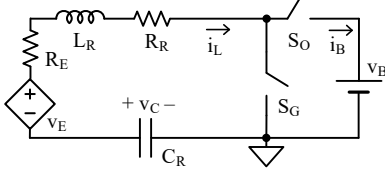


Fig. 2. Switched resonant half-bridge power stage.

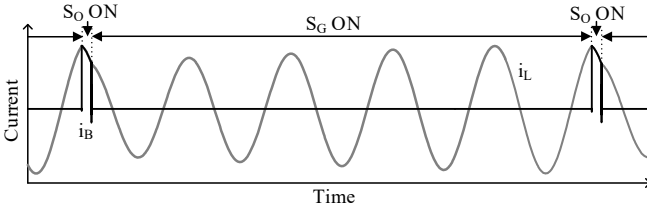


Fig. 3. Waveforms of the switched resonant half-bridge.

In a switched resonant half-bridge, the transmitter couples an open-circuit voltage  $v_E$  on the receiver coil  $L_R$  with a coupled resistance of  $R_E$  [18]. The coil  $L_R$ 's effective series resistance (ESR) is  $R_R$ .  $L_R$  and  $C_R$  in Fig. 2 resonate at  $v_E$ 's frequency. When the ground switch  $S_G$  is closed,  $v_E$  sources power into the LC tank. As  $S_G$  opens and the output switch  $S_O$  closes, the energy accumulated in the LC tank is partially drained to  $v_B$ . To raise  $i_L$  and boost power from  $v_E$ , the LC oscillation voltage  $v_C$  often exceeds the CMOS circuit's breakdown voltage  $V_{BD}$ , as  $S_G$  and  $S_O$  do not see the high voltage [10]. The over- $V_{BD}$  flexibility extends the circuit's operational coupling range for maximum power point (MPP) [9]. For MPP, the  $S_G$  and  $S_O$  need to switch synchronously near  $i_L$ 's peak, or, equivalently,  $v_C$ 's zero, as shown in Fig. 3 [18]. The power stage proposed and analyzed in [9] and [18] outputs the highest MPP for a wide coupling range. However, challenges remain in the control implementation. First, synchronization is difficult as  $v_C$  can grow beyond  $v_{BD}$ . Second, the 3-variable MPP method in [18] needs to be simplified to reduce the area and power.

This paper addresses the control implementation challenges that [18] didn't address: it presents a self-synchronization technique and a simplified MPP adjustment scheme for the switched resonant half-bridge. The rest of the paper is organized as follows. Section II discusses the design and operation of the proposed self-synchronized switched resonant half-bridge charger that uses a high-voltage-sensing comparator for synchronization. The circuit implementation is described in Section III. Section IV assesses and compares the measured performance of the fabricated prototype. Conclusions are drawn in Section V.

## II. SELF-SYNCHRONIZED SWITCHED RESONANT HALF-BRIDGE

### A. Synchronization

For MPP, the controller needs to synchronize the energy transfer with  $i_L$ 's peak [18]. Sensing  $i_L$  directly is difficult and adding sensing resistance significantly lowers the available power from the receiver coil. The other option is to sense  $v_C$ , as  $i_L$  peaks when  $v_C$  crosses zero.

However, sensing  $v_C$  is also challenging, as  $v_C$  swings between positive and negative, and its magnitude can exceed the circuit's breakdown. The high voltage needs to be divided for the circuit to sense. Although the resistive divider in Fig. 4(a) can lower the voltage, its dividing ratio is fixed. As the coupling factor  $k_C$  varies for orders of magnitude in practical applications, the divided output  $v_{SEN}$  scales proportionally. As a result of the fixed dividing ratio,  $v_{SEN}$  is too high for breakdown at high  $k_C$ , but too low for the comparator to sense at low  $k_C$ .

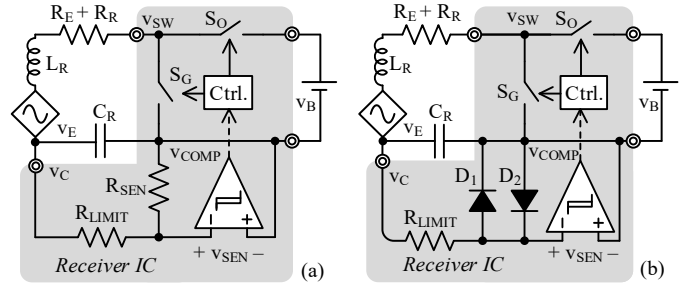


Fig. 4. Sensing high  $v_C$  with (a) resistive divider, and (b) a variable divider.

To reduce the coupling sensitivity, the diode-clamped voltage divider in Fig. 4(b) divides the voltage with a variable ratio. The circuit is composed of a pair of diodes  $D_1$  and  $D_2$  and a current limiting resistor  $R_{LIMIT}$ . When  $v_C$ 's amplitude is within a diode voltage  $v_D$ , neither  $D_1$  nor  $D_2$  conducts, so  $v_{SEN}$  follows  $v_C$ . As  $v_C$ 's amplitude grows beyond  $v_D$ , either  $D_1$  or  $D_2$  conducts current and clamps  $v_{SEN}$  at  $\pm v_D$ , as Fig. 4(b) shows. In measurement, the coupling  $k_C$  grows from 0.067% to 7.9%, so  $v_C$  varies  $36\times$  from 0.56 V to 20 V. However, as Fig. 5 shows, the divided voltage  $v_{SEN}$  varies less than  $2.7\times$  in simulation due to the  $D_1$  and  $D_2$ 's voltage suppression. This way, the dividing ratio is low ( $\ll 1$ ) at high  $v_C$ , but high ( $=1$ ) at low  $v_C$ .

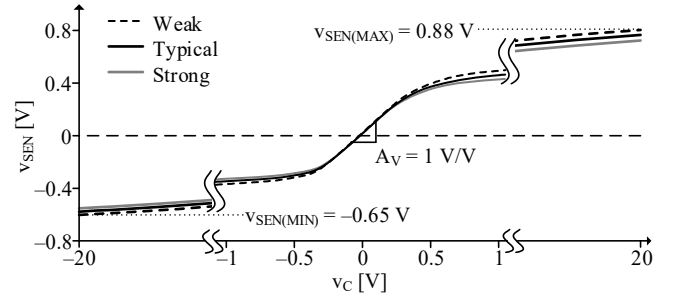


Fig. 5. Simulated output curve of the variable divider.

### B. Maximum Power Point

The goal of the inductively coupled wireless charger design is to maximize its output power to charge up the battery as fast as possible. For MPP, the receiver needs to drain just enough energy such that the averaged  $v_{C(PK)}$  over cycles is at its optimal level  $v_{C(OPT)}$  [10], where

$$v_{C(OPT)} = \left( \frac{\pi L_R f_0}{R_E + R_R} \right) v_{E(PK)}. \quad (1)$$

To maintain the  $v_{C(PK)}$  around  $v_{C(OPT)}$ , the energy transfer duration  $t_{ON}$  and energy transfer frequency  $f_X$  can be adjusted. Although adjusting both  $t_{ON}$  and  $f_X$  gives the highest power, it complicates the controller design. Fortunately, near MPP,  $P_O$  is not sensitive to  $t_{ON}$  variation. In [18],  $P_{O(MPP)}$  is lowered by less than 1.3% even if  $t_{ON}$  is 24% off its optimal value. Adjusting  $f_X$  alone gives about the same  $P_{O(MPP)}$  while simplifying control.

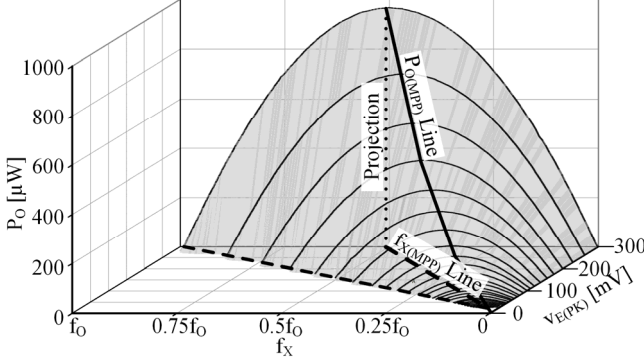


Fig. 6. Measured  $P_O$  across  $v_{E(PK)}$  and  $f_X$ .

Fig. 6 shows how  $P_O$  varies across  $f_X$  at different  $v_{E(PK)}$ . At very low  $f_X$ , the battery draws little power from the LC, so the energy in the LC builds up high. As a result, the quadratically growing ohmic loss dominates so the output power is low. Similarly, at high  $f_X$ , the battery draws too much power from the LC, so little energy remains in the LC tank. As a result,  $v_C$  stays below  $v_{C(OPT)}$ .  $P_O$  maximizes at  $f_{X(MPP)}$  when  $v_C$  averages  $v_{C(OPT)}$ . The MPP theory in [18] predicts that  $f_{X(MPP)}$  grows proportionally with  $v_{E(PK)}$ :

$$f_{X(MPP)} = \left( \frac{\pi}{4} \right) \left( \frac{v_{E(PK)}}{v_B} \right) \left/ \sin \left[ \left( \frac{t_{ON}}{t_0} \right) \pi \right] \right. \propto v_{E(PK)}. \quad (2)$$

### C. Full System

The switched resonant half-bridge in Fig. 7 uses a resonant tank  $L_R$ - $C_R$  to boost current and power from  $v_E$ . The  $L_R$  and  $C_R$  are tuned to  $v_E$ 's frequency  $f_0$ , so  $v_E$  constantly sources power into the LC tank.  $R_E$  represents the coupled resistance from the transmitter, while  $R_R$  is the ESR of  $L_R$  at  $f_0$ .

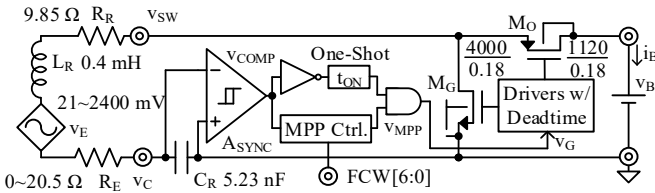


Fig. 7. Switched resonant half-bridge wireless charger.

To transfer energy, the power stage alternately switches between two modes: receiving energy from  $v_E$  and transferring received energy to  $v_B$ . The ground switch  $M_G$  closes for most of the cycle  $t_X$ , so  $L_R$ - $C_R$  receives and stores energy from  $v_E$ . Then,  $M_G$  is turned off and the output switch  $M_O$  is turned on for  $t_{ON}$ , so the energy accumulated is transferred to the battery  $v_B$ . The deadtime logic inserts around 1 ns delay so  $M_G$  or  $M_O$  only turns on when the other switch is completely off. This prevents both switches from turning on at the same time and

discharge the battery. The ground switch  $M_G$  and the output switch  $M_O$  are 4000  $\mu\text{m}$  and 1120  $\mu\text{m}$  wide, respectively. The sizes are optimized to minimize the losses for 300  $\mu\text{W}$ , which is the most probable power level for targeted glucose and blood-pressure sensing applications [2, 19].

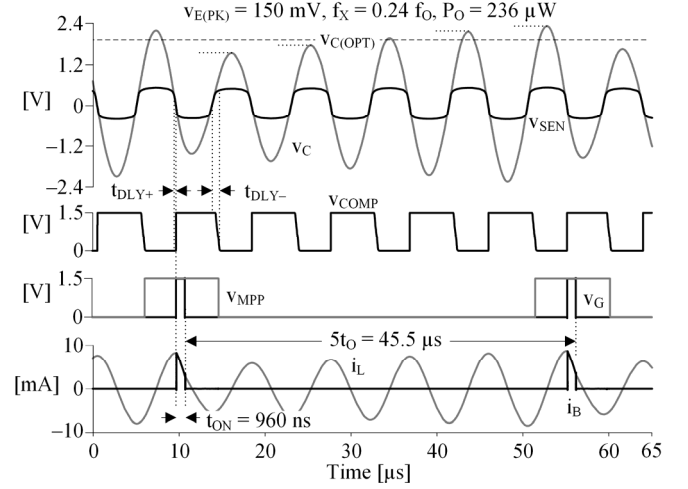


Fig. 8. Simulated wireless charger waveforms at 150 mV  $v_{E(PK)}$ .

For MPP, the energy transfer needs to be synchronized with the oscillation. For this, the synchronizing comparator  $A_{SYNC}$  in Fig. 7 detects the zero crossings of  $v_C$  in the negative direction. To transfer power near  $i_L$ 's peak, the comparator is designed to optimize its rising edge delay  $t_{DLY+}$  over its edge delay  $t_{DLY-}$ , so  $t_{DLY+} \ll t_{DLY-}$ . Once a crossing is detected, a one-shot circuit triggers a fixed pulse  $t_{ON}$ , as Fig. 8 shows. The charger only transfers power to the battery when the gating signal  $v_{MPP}$  is high. The MPP controller uses a 7-bit frequency control word  $FCW[6:0]$  to adjust the average number of cycles between energy transfers and in a delta-sigma fashion. The detailed operation will be discussed in the next section.

## III. CIRCUIT IMPLEMENTATION

### A. Synchronizing Comparator

To protect the sensing circuit from breakdown, a voltage divider lowers  $v_C$  to  $v_{SEN}$ , as shown in Fig. 9. Diode-connected NMOS  $M_{D1}$  and  $M_{D2}$  in Fig. 9 replace the diodes  $D_1$  and  $D_2$  in Fig. 4(b). The comparator in Fig. 9 compares  $v_{SEN}$  with the ground to detect  $v_C$ 's zero crossings and synchronizes the energy transfer. The comparator needs to (i) take negative input as  $v_{DIV}$  swings from  $-v_D$  to  $+v_D$ , and (ii) minimize the  $v_C$ 's falling edge delay  $t_{DLY-}$  for MPP. For (i), a PMOS pair,  $M_1$  and  $M_2$ , is used.  $M_1$ 's gate is grounded which generates a bias voltage  $v_B$  for  $M_4$ 's pull-up path. Above-zero  $v_{DIV}$  crushes  $M_4$ 's  $v_{GS}$ , so  $M_6$  pulls  $v_{O1}$  up slowly with fixed 180 nA. As  $v_{DIV}$  drops below zero,  $M_2$ 's current grows quadratically with the voltage drop and pulls  $v_{O1}$  down quickly. As  $v_{DIV}$  drops below zero,  $M_2$ 's source and body follow, preventing the circuit from breakdown. For (ii), a secondary common-source stage  $M_{10}$  expedites  $v_{COMP}$ 's pulling-up as  $v_{COMP1}$  falls. Combining that  $v_{COMP1}$  pulls down fast and  $v_{COMP}$  pulls up fast, the measured  $t_{DLY-}$  (120 ns  $\sim$  180 ns) is much shorter than  $t_{DLY+}$  (1  $\mu\text{s}$   $\sim$  1.8  $\mu\text{s}$ ), as Fig. 8 shows.

The voltage divider induces loss as the  $R_{LIMIT}$  steals and burns a fraction of  $i_L$ . At high  $v_C$ , ignoring the voltage drop on the diodes, the fractional loss that parallel  $R_{LIMIT}$  induces is equivalent to a series resistance of  $0.3\Omega$  [20], which lowers  $P_{O(MPP)}$  by 3%. At low  $v_C$ , as the diodes' voltage drop lowers the current across  $R_{LIMIT}$  more,  $R_{LIMIT}$ 's loss is less than 3%. A larger  $R_{LIMIT}$  lowers the current stolen and the loss. However, larger  $R_{LIMIT}$  also adds more parasitic capacitance at  $v_{SEN}$ , causing more delay. An offset time that equals 3% of the period lowers the output power by 4% [18]. A  $250\text{ k}\Omega$   $R_{LIMIT}$  is chosen to keep the RC delay within 3% of the oscillation period.

\* NMOS bulk connected to GND. PMOS bulk connected to source

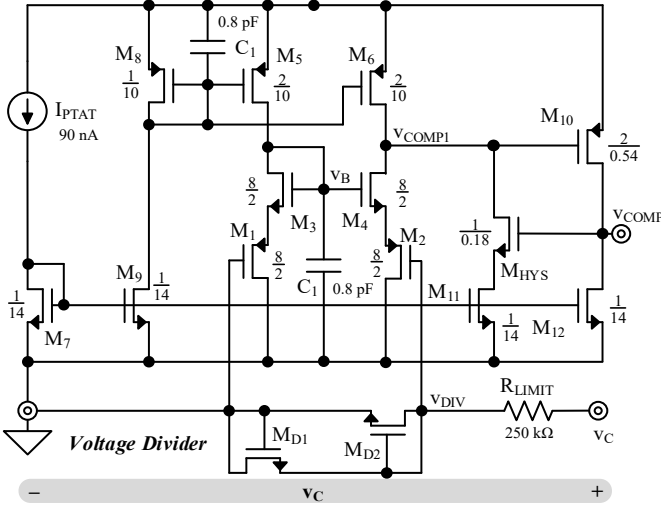


Fig. 9. Synchronizer circuit with the variable divider

### B. MPP Controller

The MPP controller is implemented with a 7-bit-input, 1-bit-output digital-to-digital delta-sigma modulator frequently found in fractional-N PLL designs [21]-[22], as shown in Fig. 10. Each time the synchronizing comparator's output  $v_{COMP}$  toggles in the negative direction, the full adder self-adds the 7-bit frequency control word FCW to the 8-bit register array ACC. This way, ACC accumulates FCW each cycle, as Fig. 11 shows. Once ACC's most significant bit (MSB),  $v_{MPP}$ , toggles high, it enables the energy transfer in the next coming cycle. The 7 least significant bits (LSBs) are fed back to the full adder's input for the next cycle's accumulation in a delta-sigma fashion. Since  $v_{ACC}$  self-accumulates  $v_{FX}$  each time and its MSB toggles once the accumulation reaches 128, the averaged energy transfer frequency is

$$f_{X(MPP)} = \frac{v_{FX}}{128} f_O \quad (3)$$

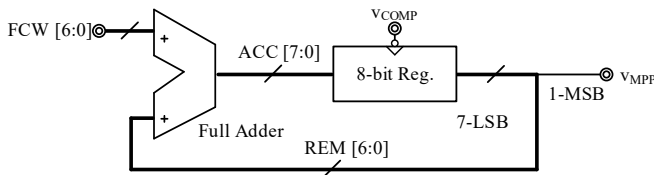


Fig. 10. Delta-sigma MPP controller.

This way,  $f_X$  can be adjusted with  $1/128 f_O$  resolution. Compared to skipping an integer number of cycles [23], the delta-sigma modulator allows skipping a fractional number of

cycles. Details including waveforms of skipping a fractional number of cycles will be discussed in the next section.

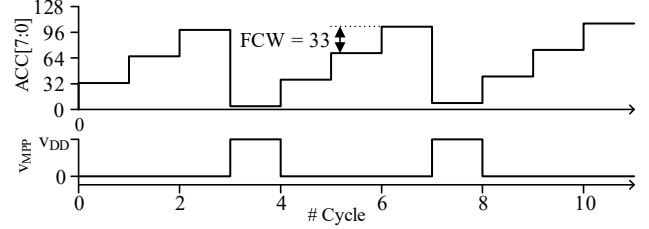


Fig. 11. Waveforms of the MPP controller.

## IV. MEASUREMENTS

### A. Prototype

To demonstrate the charger's functionality and performance, a prototype is built with 180-nm CMOS technology. The charger IC, as shown in Fig. 12, integrates the power stage ( $M_G$ ,  $M_O$ , and the driver) and the synchronizer ( $I_{PTAT}$ ,  $A_{SYNC}$ ,  $V_{DIV}$ ) while occupying only  $220\ \mu\text{m} \times 381\ \mu\text{m}$  of silicon area. The prototype operates at 110 kHz, which is much lower compared to the 6.78 MHz used in the previous design [9]. The lower frequency operation greatly reduces the power consumption of the synchronizing comparator. The downside is that the resonant capacitor  $C_R$  increases to 5.23 nF so it cannot be integrated on-chip. However, the off-chip  $C_R$  measures only  $1.6 \times 0.8 \times 1.6\ \text{mm}^3$ , which adds a limited volume to the overall system. The wireless charger uses the 0.4 mH Coilcraft 4513TC as the receiver coil that measures  $11.7 \times 3.5 \times 2.6\ \text{mm}^3$ . This is smaller than the  $4.7\ \mu\text{H}$  Coilcraft PA6512-AE used in [9]. The MPP control and one-shot circuit in Fig. 2 are implemented on an FPGA for testing flexibility. The MPP control is fully synthesizable and can be migrated on-chip easily.

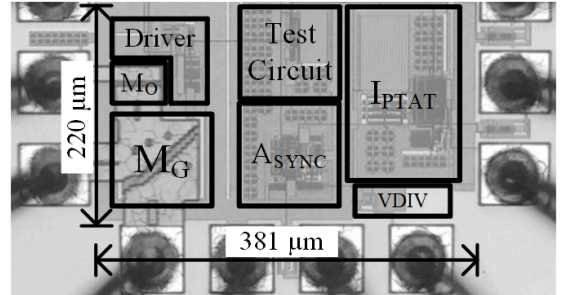


Fig. 12. Photos of the wireless charger IC.

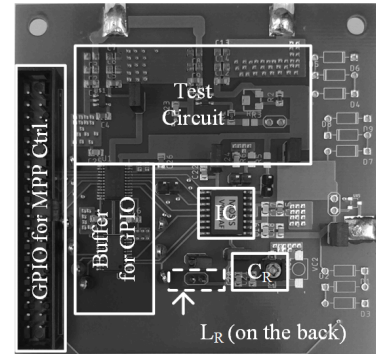
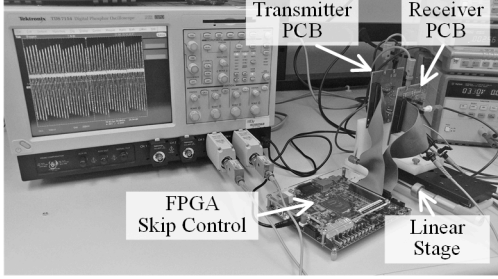
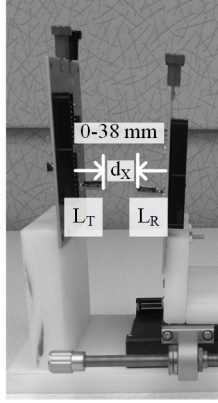


Fig. 13. Photos of the wireless charger PCB.

Fig. 13 shows the photo of the charger prototype PCB. The FPGA and the charger PCB are connected via the GPIO port, as shown in Fig. 13. The linear stage in Fig.14 adjusts the distance between the wireless charger and the source from 0 to 38 mm. As a result, the coupled open-circuit voltage  $v_E$  on the receiver coil ranges from 24 mV to 2.8 V. The transmitter couples up to  $20.5 \Omega$  back on the receiver coil.



(a) Measurement setup



(a) Linear stage that adjusts  $d_X$  from 0 to 38 mm

Fig. 14. The measurement setup and the linear stage.

Resonance boosts power but puts electrical stress on the receiver. The electrical stress grows with the oscillation until saturated by the ohmic loss [10]. When saturated, the oscillation voltage is twice as high as  $v_{C(OPT)}$  [10], so

$$v_{C(MAX)} = 2v_{C(OPT)} = \left( \frac{sL_R}{R_E + R_E} \right) v_{E(PK)}, \quad (4)$$

where  $sL_R$  is  $L_R$ 's impedance at  $f_0$ . As  $C_R$  and  $L_R$  resonate at  $f_0$ ,  $v_C + v_L \approx 0$ . Therefore,

$$i_{L(MAX)} = \frac{v_{E(PK)}}{R_E + R_R}. \quad (5)$$

With the measured  $v_{E(PK)}$  and  $R_E$ ,  $C_R$  and  $L_R$ 's voltage and current stress can be calculated using (4) and (5). In measurement,  $v_{C(MAX)}$  ranges 20~0.57 V, and  $i_{L(MAX)}$  ranges 74~2.1 mA across 0~38 mm.

On the transmitter side, a power inverter (modeled as  $v_S$  in Fig. 1) drives the  $L_T C_T$  resonant tank. As the receiver draws little power the current and voltage stress on  $L_T$  and  $C_T$  relatively constant across  $d_X$ . Transmitter's  $i_{T(MAX)}$  and  $v_{T(MAX)}$  can be calculated similarly as the receiver:

$$i_{T(MAX)} = \left( \frac{4}{\pi} \right) \left( \frac{v_{S(PK)}}{R_S + R_T} \right), \quad (6)$$

$$v_{T(MAX)} = \left( \frac{sL_T}{R_S + R_T} \right) \left( \frac{4}{\pi} \right) v_{S(PK)}, \quad (7)$$

where  $R_T$  is the source impedance of the driving inverter. Compared to (4) and (5), a  $4/\pi$  term is multiplied to  $v_{S(PK)}$  as the sinusoidal fundamental tone of the square-wave  $v_S$  is  $4/\pi$  times higher [18]. In the measurement setup,  $v_{S(PK)} = 3.3V$ ,  $R_S = 5.6 \Omega$ , and  $R_T = 10.2 \Omega$ . The derived  $v_{T(MAX)}$  and  $i_{T(MAX)}$  are 73 V and 265 mA, respectively.

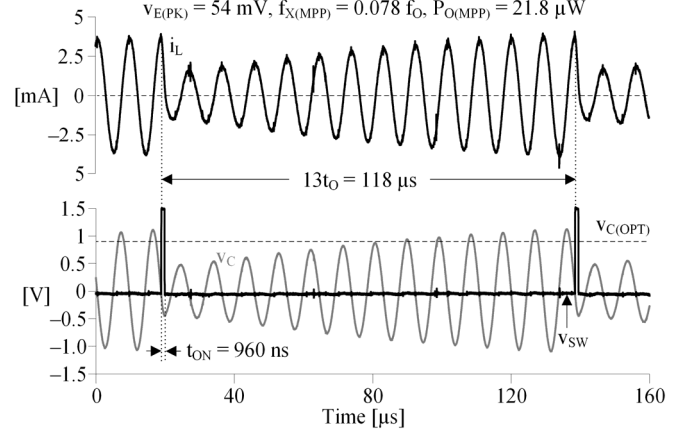


Fig. 15. Measured MPP receiver waveforms when  $k_C = 0.15\%$ .

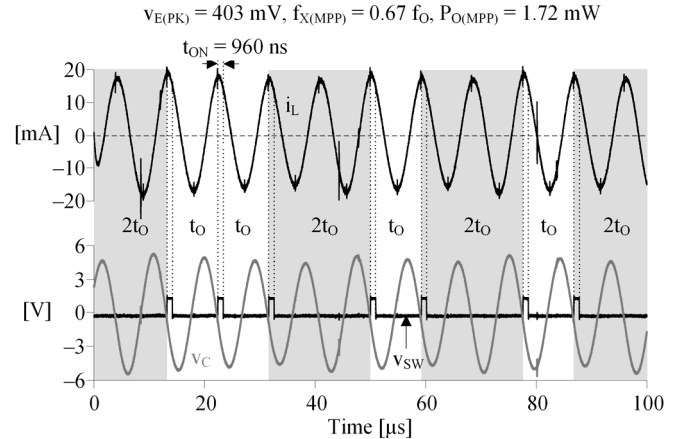


Fig. 16. Measured MPP receiver waveforms when  $k_C = 1.1\%$ .

Fig. 15 and Fig. 16 show the measured  $v_{SW}$ ,  $v_C$ , and  $i_L$  waveforms at MPP when the receiver is 28 mm and 10 mm away from the power source. The couplings are 0.15% and 1.1%, respectively. When the coupling is weak, the charger transfer energy less frequently to allow energy to build up in the LC tank and raise  $L_R$ 's voltage. At 0.15%, the wireless charger on average skips 12.8 cycles between energy transfer for MPP. In other words,  $f_{X(MPP)} = 0.078 f_0$ . When the coupling is high, the charger transfer energy more frequently to avoid LC energy build up too high and causes excessive conduction loss on  $L_R$ . At 1.1% coupling, the wireless charger transfers energy every one or two cycles, such that the average number of cycles between energy transfer is 1.5. So the effective MPP transfer frequency  $f_{X(MPP)} = 0.67 f_0$ .

### B. Charging Profile

To evaluate the charger's performance, a charging test is performed that charges up a 1.1  $\mu F$  capacitor  $C_O$  from 1 V to 1.8 V when the receiver is 13 mm, 18 mm, 23 mm, and 28 mm away from the power source. At the above distances, the

receiver fully charges  $C_O$  in 3.6 ms, 7.9 ms, 20 ms, and 63 ms, respectively. The corresponding charging currents are 500  $\mu\text{A}$ , 250  $\mu\text{A}$ , 9.9  $\mu\text{A}$ , 3.2  $\mu\text{A}$ , respectively, as shown in Fig. 17.

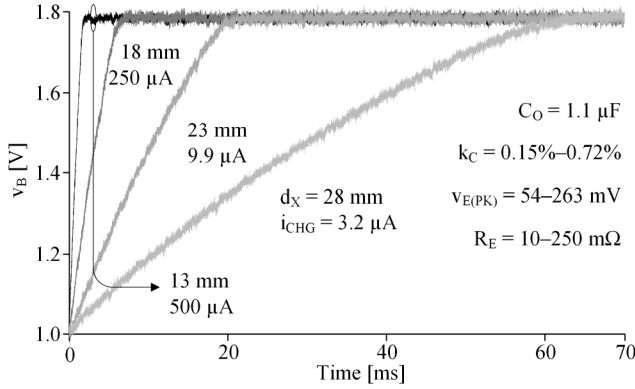


Fig. 17. Measured charging waveforms at 13~28 mm distance.

### C. Ideality Factor

The maximum output power  $P_{O(MPP)}$  is the key performance of the wireless charger. However, different wireless chargers'  $P_{O(MPP)}$  is not comparable, as  $P_{O(MPP)}$  scales with the transmitter's power and coupling. To assess the relative performance of the wireless charger,  $P_{O(MPP)}$  needs to be normalized. Maximum available power  $P_{O(MAX)}$  defines the highest power the receiver can possibly draw from the transmitter at the given coupling:

$$P_{O(MAX)} = \left( \frac{0.5v_{E(PK)}}{\sqrt{2}} \right)^2 \left( \frac{1}{R_R + R_E} \right), \quad (8)$$

$P_{O(MAX)}$  typically drops with  $d_X^6$  when distantly coupled [24]. The ideality factor  $\eta_I$  references  $P_{O(MPP)}$  to  $P_{O(MAX)}$  and normalizes  $P_{O(MPP)}$  with non-receiver variables:

$$\eta_I = \frac{P_{O(MPP)}}{P_{O(MAX)}}. \quad (9)$$

Fig. 18 shows the measured  $\eta_I$  when  $d_X$  ranges from 0–38 mm. Ideality  $\eta_I$  is high ( $> 60\%$ ) at 7–27 mm. Past 27 mm,  $\eta_I$  gradually drops to zero. This is because, as the power source separates further from the wireless charger, it couples less  $v_E$  on the receiving coil and, according to (2), the wireless charger skips more cycles between energy transfer for MPP. Although on average,  $v_C$  is at  $v_{C(OPT)}$ , each cycle's  $v_{C(PK)}$  deviates further from  $v_{C(OPT)}$ , the non-linear loss [18] caused by the deviation grows and lowers  $\eta_I$ .

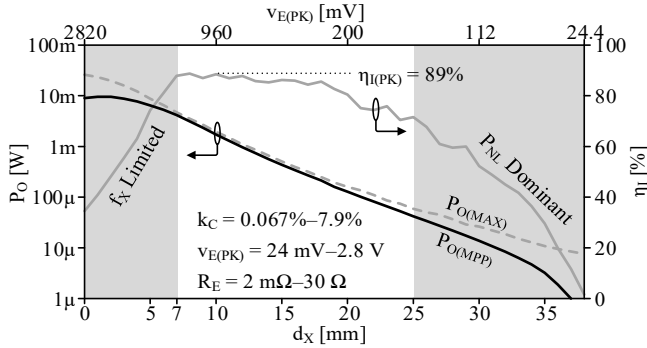


Fig. 18. Measured ideality factor across power transmission distance  $d_X$ .

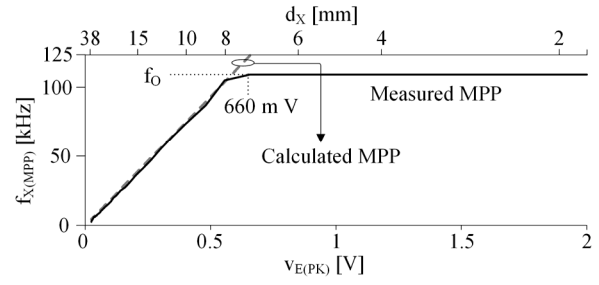


Fig. 19. Measured  $f_{X(MPP)}$  across  $v_{E(PK)}$ .

On the other end, when  $d_X$  is shorter than 7 mm,  $\eta_I$  is also low. This is because, as the transmitting source couples more  $v_E$  on the receiving coil,  $f_{X(MPP)}$  scales up linearly and eventually reaches its maximum limit:  $f_o$ . Beyond that, the MPP controller can no longer adjust  $f_X$  to the desired  $f_{X(MPP)}$ , so  $\eta_I$  starts to drop. As shown in Fig. 19,  $f_{X(MPP)}$  reaches  $f_o$  as  $v_{E(PK)}$  grows above 0.58 V. Then  $f_{X(MPP)}$  is capped to  $f_o$ , and  $P_{O(MPP)}$  becomes  $f_X$ -limited. Transferring energy for a longer duration  $t_{ON}$  improves  $\eta_I$  when the coupling is high and  $P_{O(MPP)}$  is  $f_X$ -limited. However, for targeted biomedical implant applications, high coupling is unlikely. Plus, adjustable  $t_{ON}$  requires additional circuitry and quiescent power. So  $t_{ON}$  is fixed at 960 ns, which is the optimal  $t_{ON}$  when coupling  $k_C$  is halfway (in log scale) across the 0.067%–7.9% range.

Fig. 20 shows the loss breakdown of the proposed inductively coupled wireless charger. Among all the losses, the nonlinear loss  $P_{NL}$  dominates when  $k_C$  is lower than 0.3% or higher than 2%. At low  $k_C$ , the power receiver skips more cycles for MPP [18], so  $P_{NL}$  is high. At high  $k_C$ ,  $P_{O(MPP)}$  becomes  $f_X$ -limited, so  $P_{NL}$  is also high. The high-voltage-sensing circuit consumes  $P_{R(SEN)}$ , as  $R_{LIMIT}$  steers away a fraction of the LC tank's current. However, when  $v_C$  drops below a diode voltage, the diodes  $D_1$  and  $D_2$  conduct little current, so  $P_{R(SEN)}$  gradually drops to zero at low  $k_C$ . The quiescent loss  $P_Q$  and charge loss  $P_C$  do not scale proportionally with  $k_C$ . Therefore, their portion grows at low  $k_C$ .

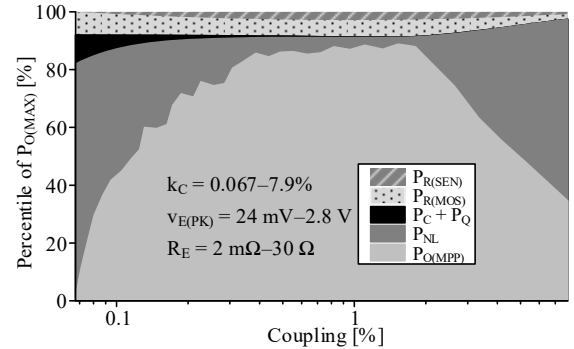


Fig. 20. Measured  $f_{X(MPP)}$  across  $v_{E(PK)}$ .

The power and ideality in Fig. 19 are measured when the transmitter coil  $L_T$  and the receiver coil  $L_R$  align in the center. In real applications, the coils' centers often misalign by a lateral distance  $d_Y$ . The math model in [25] predicts the magnetic field  $H_Z$  that  $L_R$  captures at distance  $d_X$  with lateral misalignment  $d_Y$ :

$$H_Z = \left( \frac{I_T}{2\pi d_Y} \right) \sqrt{\frac{m}{4r_T d_Y}} \left[ d_Y K + \frac{r_T m - (2-m)d_Y}{2-2m} E \right], \quad (10)$$

TABLE I: RELATIVE PERFORMANCE

	JSSC '13 [16]	TCAS II '13 [17]	JESTPE '13 [18]	AICASP '20 [3]	TCAS I '16 [30]	TCAS II '18 [9]	JSSC '16 [31]	This Work
<b>Receiver Type</b>	Switched Bridge				Switched Resonant Half-Bridge			
<b>Tech. (<math>\mu\text{m}</math>)</b>	0.18	0.18	0.18	0.5	Board	0.18	0.18	0.18
<b>Silicon Area (<math>\text{mm}^2</math>)</b>	0.49	0.245	0.26	2.25	–	0.471	0.544	0.084
<b>Tx Coil Size</b>	224 $\text{mm}^3$	224 $\text{mm}^3$	–	100 $\text{mm}^3$	28900 $\text{mm}^2$	633 $\text{mm}^3$	–	106 $\text{mm}^3$
<b>Rx Coil Size</b>	106 $\text{mm}^3$	106 $\text{mm}^3$	106 $\text{mm}^3$	25 $\text{mm}^3$	900 $\text{mm}^2$	633 $\text{mm}^3$	106 $\text{mm}^3$	106 $\text{mm}^3$
<b><math>d_x</math> (mm)</b>	0–11	0–11	10–50	0–8	70	13–38	85	0–38
<b><math>f_0</math> (MHz)</b>	0.125	0.125	0.125	13.56	1	6.78	0.05	0.11
<b><math>k_C</math> (%)</b>	0.59–6.7	0.9–7.6	0.15–1.35	–	1.3	0.09–1.1	–	0.067–7.9
<b><math>k_C</math> Attenuation w/ <math>d_Y</math> (dB)</b>	–	–	–	–	–	–	–	–20 @ $d_X = d_Y$
<b><math>V_{E(PK)}</math> (mV)</b>	39.5–386	46–480	66–585	–	41	18.5–282	–	24–2880
<b><math>P_{O(MPP)}</math> (<math>\mu\text{W}</math>)</b>	0–224	26.6–830	16–557	5830	96.1	1.2–1340	0–2.84	0.1–26000
<b>Over-<math>V_{BD}</math> Operation</b>	No	No	No	No	No	Yes	No	Yes
<b>Self-Synchronized</b>	No	No	Yes	Yes	No	No	Yes	Yes
<b>Closed-Loop MPPT</b>	No	No	No	Yes	No	No	Yes	No
<b><math>\eta_{I(PK)}</math> (%)</b>	30.9*	46.9*	28.6*	–	42.9	84.8	67.7	88.8
<b><math>k_C</math> Range</b>	9.8 $\times$	8.4 $\times$	9 $\times$	–	–	12.2 $\times$	7.1 $\times$	118 $\times$

\*Ideality inferred from the reported  $P_{O(MPP)}$ ,  $V_{E(PK)}$ ,  $R_R$ , and estimated  $R_E$ , using (3) and (4).  $R_E$  is estimated as in [12].

$$m = \frac{4r_T d_Y}{(r_T + d_Y)^2 + d_X^2} \quad (11)$$

Here,  $r_T$  is the radius of the transmitter coil.  $K$  and  $E$  are the complete elliptic integrals of the first and second kind [25]. Lateral misalignment attenuates  $H_Z$  and thus lowers the coupling  $k_C$ . Fig. 21 shows how  $k_C$  attenuates with  $d_Y$  at 10, 20, and 30-mm separation using Eqn. (10)'s prediction. At  $d_Y = d_X$ ,  $V_{E(PK)}$  attenuates  $\sim 20\text{dB}$  across different  $d_X$ . Ideality  $\eta_{I(PK)}$  for misaligned coils can be estimated using Fig. 20 and Fig. 21. First, use Fig. 21 to estimate the attenuated  $k_C$  at  $d_Y$ . Then, map the attenuated  $k_C$  onto Fig. 20 to find the estimated  $\eta_{I(PK)}$ .

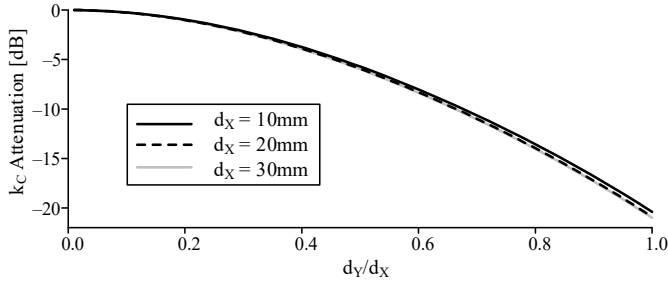


Fig. 21. Measured  $k_C$  across  $V_{E(PK)}$ .

#### D. Relative Performance

Fig. 22 compares the proposed charger's  $\eta_I$  across  $k_C$  with the state of the art. For a fair comparison, single-stage wireless rectifying-regulating supply systems [14, 26–29] are excluded, as their design goal is to maximize the PCE of the receiver instead of MPP. The wireless chargers in [30], [9], and [31] are based on switched resonant half-bridge and its variations. Although [9] achieves an ideality as high as 85%, the charger cannot self-synchronize. Therefore, the system is incomplete. The charger in [30] has no synchronizer either. Plus, as the charger completely drains the LC tank's energy each time, the energy transfer is not optimal. So ideality is low at 42.9%. The charger in [31] is self-synchronized and achieves 67.7% of  $\eta_I$ . However, its power stage cannot operate beyond the circuit's

breakdown, so the coupling range is  $16\times$  narrower. The charger in [31] includes a maximum power point tracking (MPPT) that this design does not have. However, the MPPT is based on a one-time calibration, so it does not affect the ideality. [31] is not included in Fig. 22, as the coupling information cannot be extracted from the paper.

The switched bridges in [15], [16], and [17] energizes and de-energizes  $L_R$  directly with the battery. As the switches in the switched bridge see the inductor voltage, the circuit cannot operate beyond breakdown. So the workable coupling ranges are much lower compared to the proposed design. The chargers in [15] and [16] cannot self-synchronize, so the systems are incomplete. Although the charger in [17] includes an integrated synchronizer, the synchronizer needs to break the charging operation, resulting in additional opportunity loss. As a result, its  $\eta_{I(PK)}$  is the lowest at 28.6%.

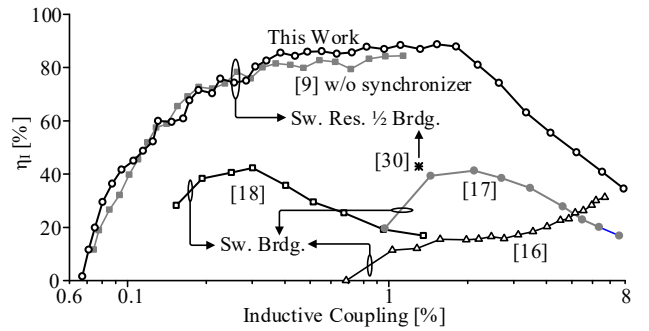


Fig. 22. Measured  $\eta_I$  compared with the state of the arts.

Table I summarizes and compares the performance of the proposed design with the state of the arts. Although the design in [3] is a supply, it adjusts the transmitter power to deliver the right amount of power to the load. In other words, it achieves MPP the same way as a charger and is thus comparable. The chargers in [30], [9], [15], and [16] are incomplete as they do not include a synchronizer. Still, among all designs, the proposed design achieves the highest  $\eta_I$  of 89% and the widest  $k_C$  range of  $118\times$ . Compared to other self-synchronized

wireless chargers, the proposed wireless charger improves  $\eta_i$  and  $k_C$  range by  $1.3\times$  and  $13\times$ , respectively.

## V. CONCLUSIONS

Embedded microsensors are critical blocks in the biomedical field as well as in the IoT. Inductively coupled battery charger greatly extends the microsensors' lifetime. For practical applications, the charger needs to output the highest power possible over a wide coupling range. This paper proposes a self-synchronized switched resonant half-bridge inductively coupled battery charger. The charger's power stage's series resonant capacitor withstands the high coil voltage needed to boost output power. The proposed charger also adjusts energy transfer patterns so the output power can stay at MPP as coupling varies orders of magnitude. Finally, the charger addresses the challenge of synchronization, as the voltage being sensed for synchronization can far exceed the CMOS circuit's breakdown. A prototype charger is fabricated using 180-nm CMOS technology. Measurements show that the proposed charger improves the output power and workable coupling range by  $1.3\times$  and  $13\times$  over the state of the art.

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